



LIII Congresso SAIt
Pisa, 4-8 Maggio 2009
L'Universo quattro secoli dopo Galileo

AstroFisica
A small map of Italy is shown with a red star indicating the location of Pisa.



Improved results on extraction of
 $^{11}\text{B}(\text{p},\alpha_0)^8\text{Be}$ S(E)-factor through the
Trojan Horse Method

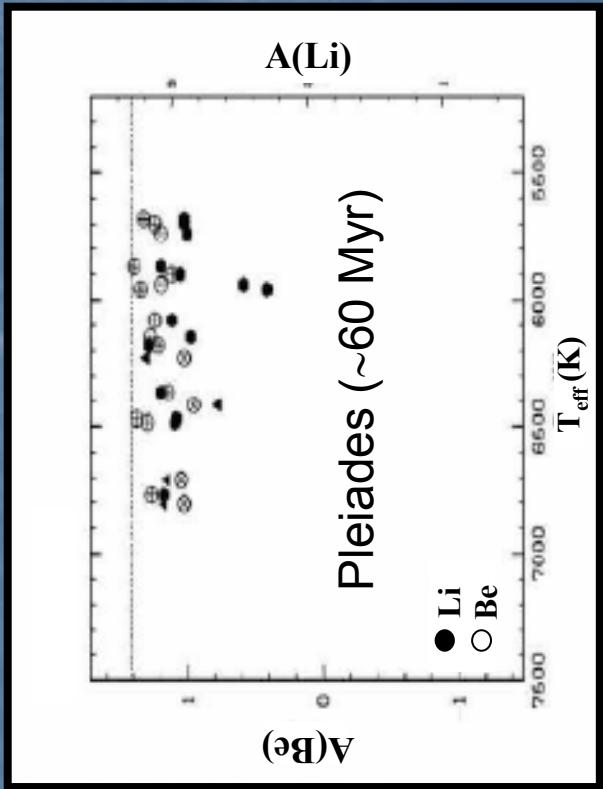
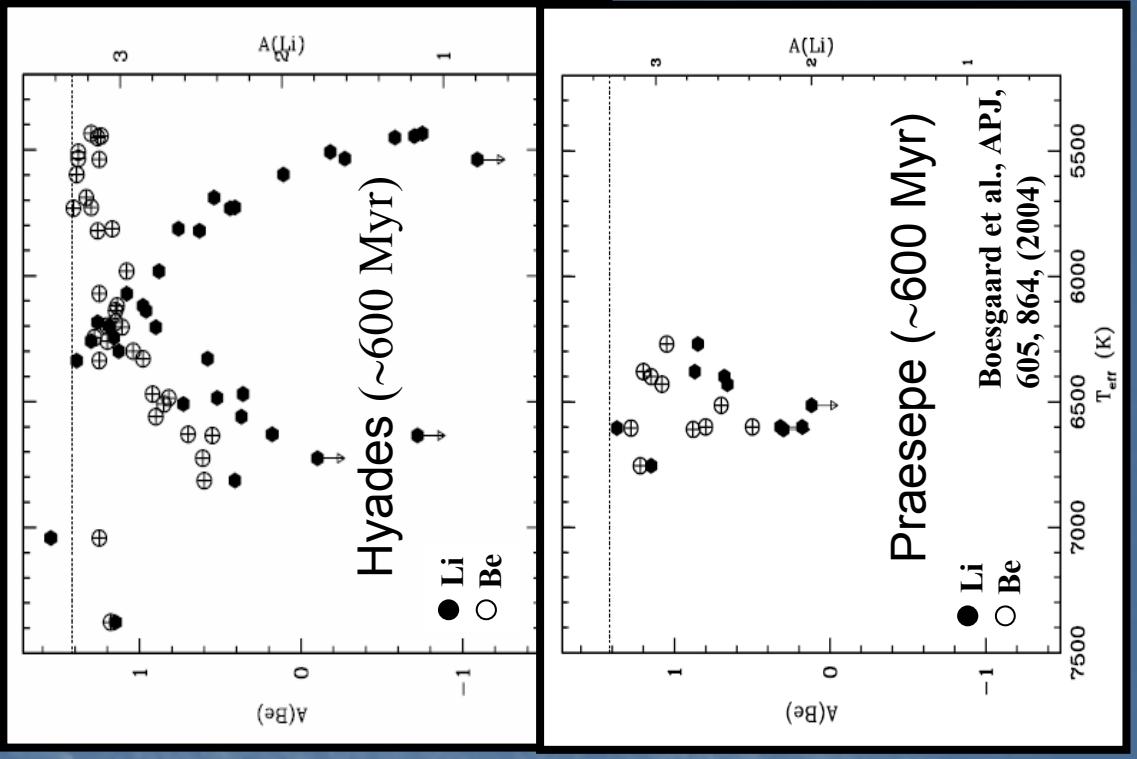


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Light Elements in Astrophysics: Lithium, Beryllium, Boron (I)

- ✓ Observational status: Depletion in Open Clusters for stars with $5500 < \text{Teff(K)} < 7500$.
- ✓ Burning (p,α) channel as the main contribution to their destruction at $T_6 = 2.5$ (Li), $T_6 = 4$ (Be), $T_6 = 5$ (B)
- ✓ Dip di Li & Be: the depth of the dip reflects the *nuclear fate* in the nuclear destruction zone (NDZ).
- ✓ Li-Be & Be-B correlation as signature of “rotation-induced” slow-mixing processes.



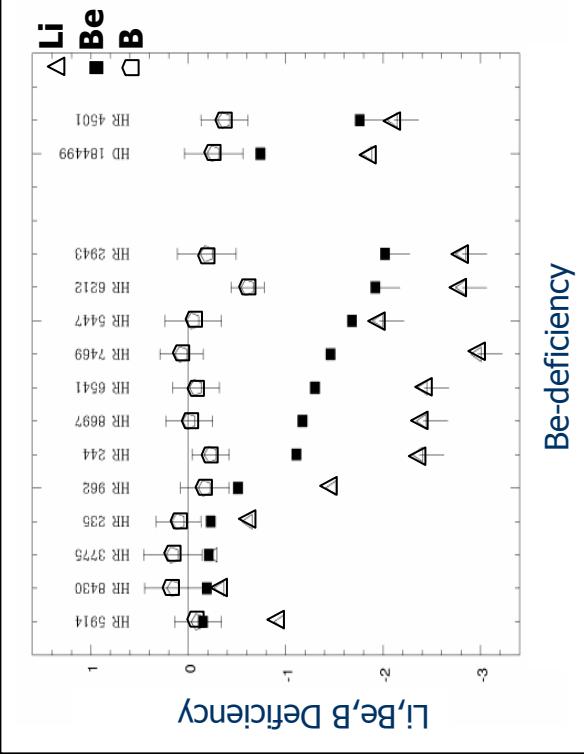
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Light Elements in Astrophysics: Lithium, Beryllium, Boron (II)

Nuclear

✓ Observational Status

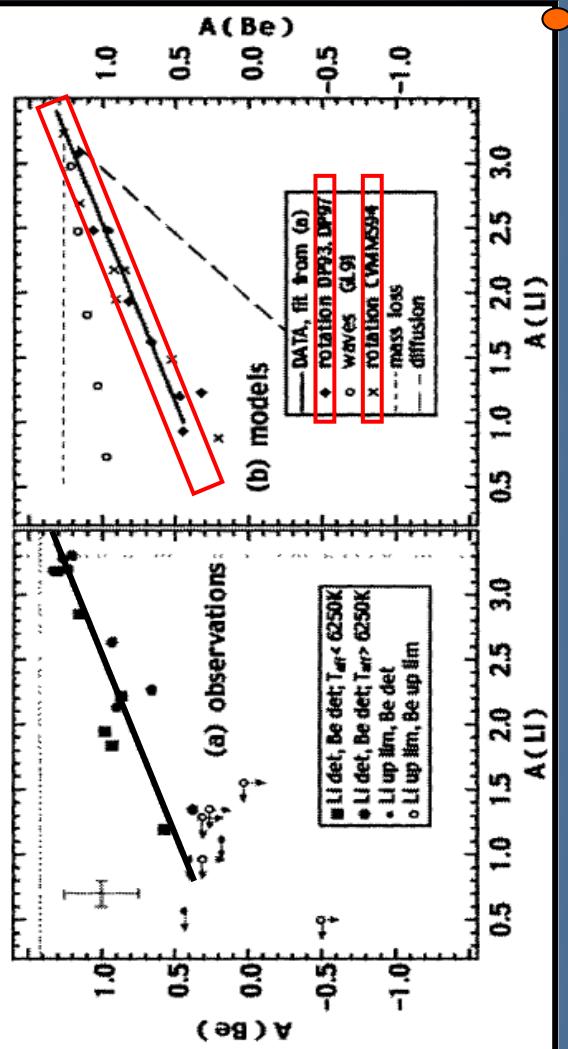
Light elements Li, Be and B abundances are shown as the Be-deficiency varies for low-metallicity young star ($5800 < \text{Teff} (\text{K}) < 6500$, $M_* \sim 1 M_\odot$).



Delyannis C.P., ASP Conference series, Vol. 198, 2000

✓ Theoretical Models

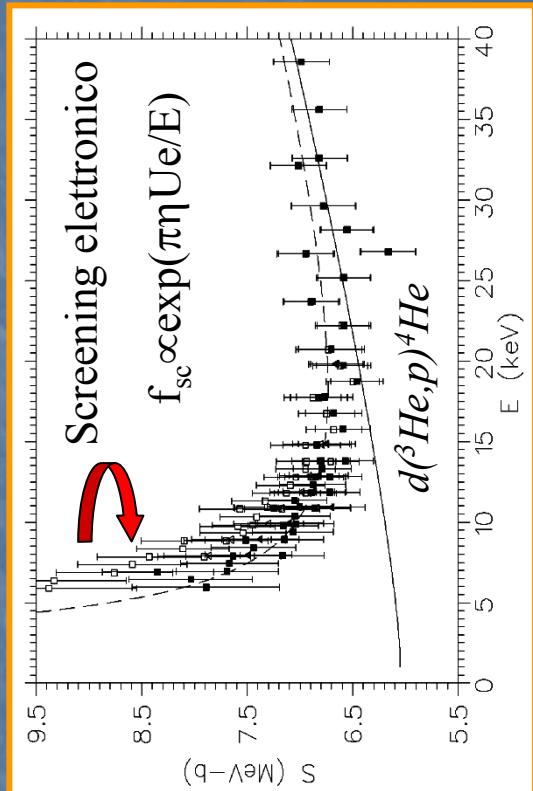
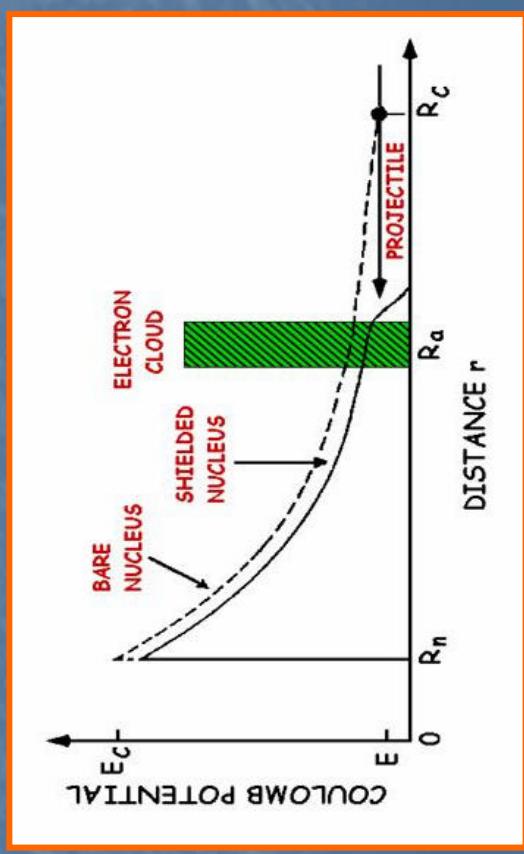
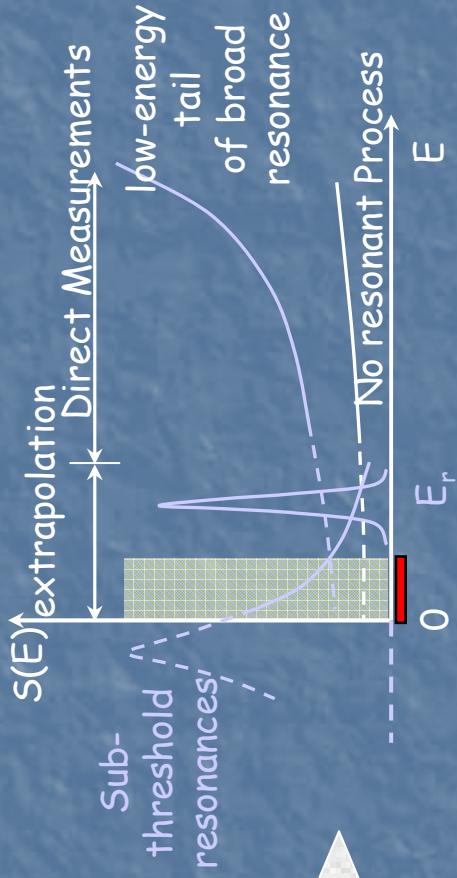
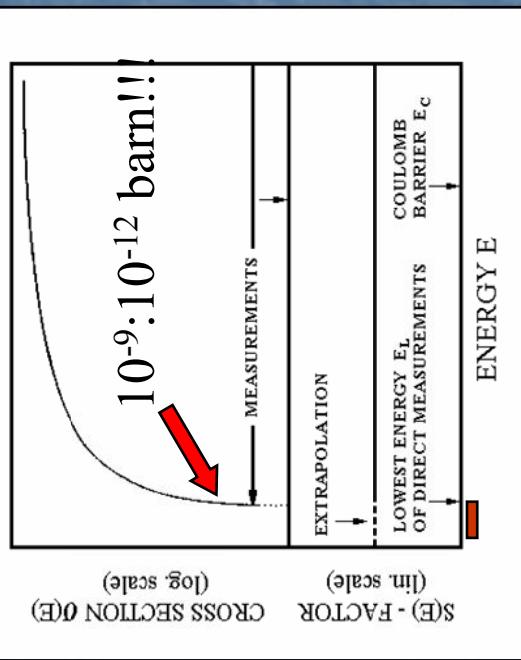
The lithium and beryllium deficiency and the boron observational trend are in agreement with “non-standard” mixing process (*slow-mixing process*) in stellar interior mainly due to rotation (Stephens et al. 1997; Boesgaard et al. 1997).



Direct Measurements in Nuclear Astrophysics and related difficulties

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Nuclear



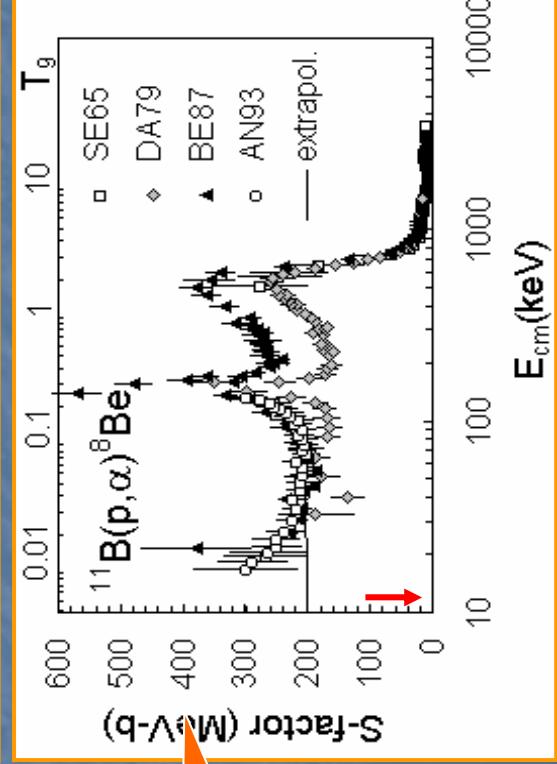
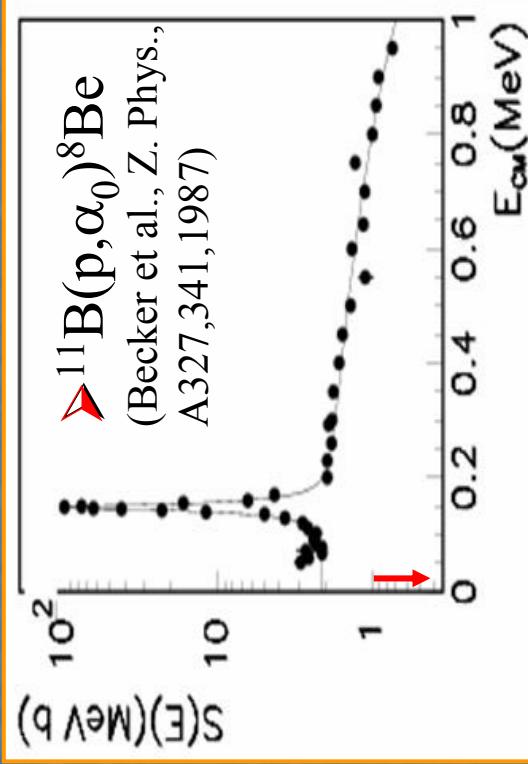
From Stars to the laboratory: direct measurements of $^{11}\text{B}(\text{p},\alpha)^8\text{Be}$ reaction

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Nuclear

- ✓ By considering the typical temperatures of some 10^6 K at which burning (p,α) reactions typically occur for boron isotopes in stellar environments, the Gamow peak is at about

$$E_0 = 1.22(Z_x^2 Z_y^2 \mu T_6^2)^{1/3} \text{ keV} \approx 10 \text{ keV}$$



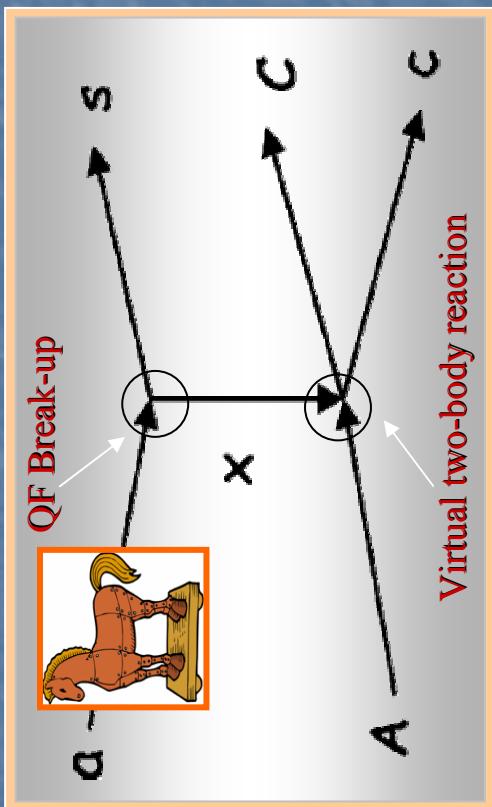
Only extrapolations are possible for low energies because screening and Coulomb penetration effects

The Indirect Methods: Trojan Horse Method



Nucleare

- ✓ The THM allows to extract a charged particles two-body cross section $x(A,C)c$ at astrophysical energies by selecting the quasi-free (QF) contribution of a suitable three body reaction $a(A,Cc)s$ performed at energies well above the Coulomb barrier.



$$\frac{d^3 \sigma}{dE_c d\Omega_c d\Omega_C} \propto KF |\Phi(p_S)|^2 \left(\frac{d\sigma}{d\Omega} \right)^N$$

(I) 2-body data free of Coulomb suppression and electron screening effects;

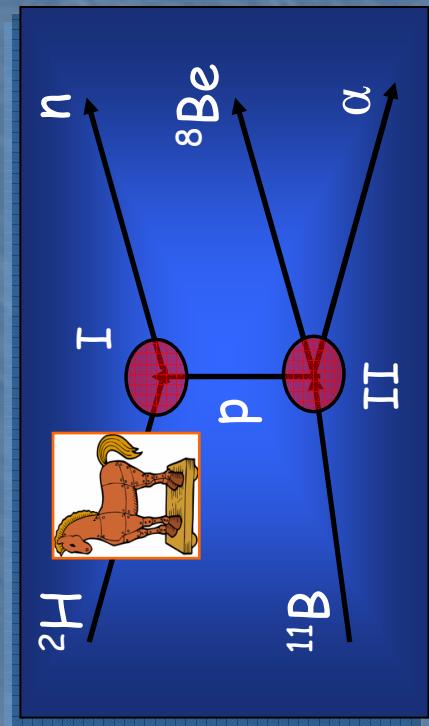
(II) validity test, introduction of penetrability function and normalization to direct data are needed!!

The explored energy region E_{cm} goes from $0 < E_{cm} < 1$ MeV by using only one value for the energy beam!!

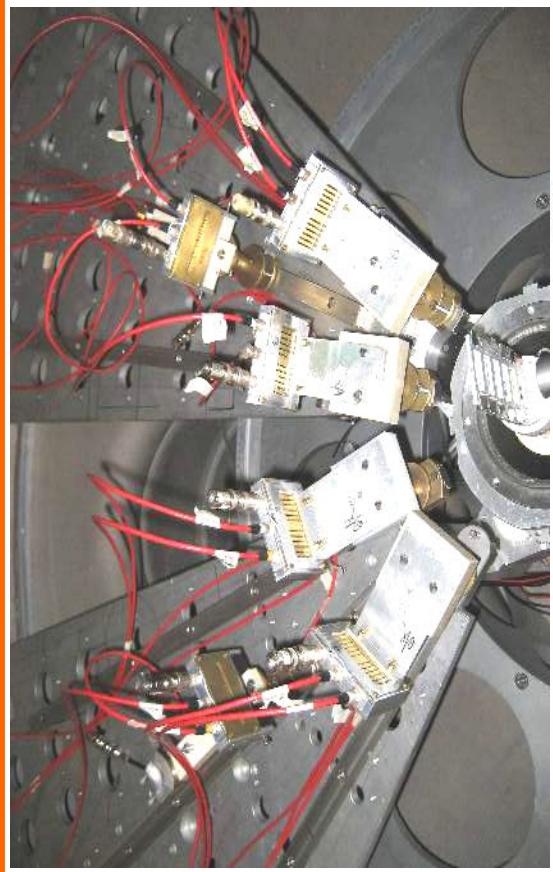


no extrapolation for $A(x,c)C$ in the energy window relevant for astrophysics

Indirect study of the $^{11}\text{B}(\text{p},\alpha_0)^8\text{Be}$ reaction through $^2\text{H}(^{11}\text{B},\alpha_0^8\text{Be})\text{n}$ by means of the THM



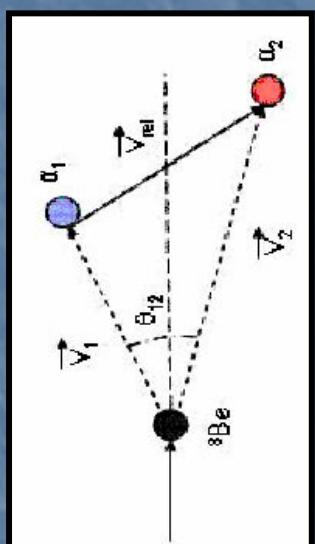
- ✓ Study of the $^{11}\text{B}(\text{p},\alpha)^8\text{Be}$ reaction ($Q=8.59 \text{ MeV}$) through $^2\text{H}(^{11}\text{B},\alpha^8\text{Be})\text{n}$ 3-body reaction ($Q=6.36 \text{ MeV}$, $E_{\text{coul}}=1.3 \text{ MeV}$);
- ✓ The deuteron is used like Trojan Horse nucleus. The momentum distribution of intercluster motion inside the deuteron is known from independent experiments.



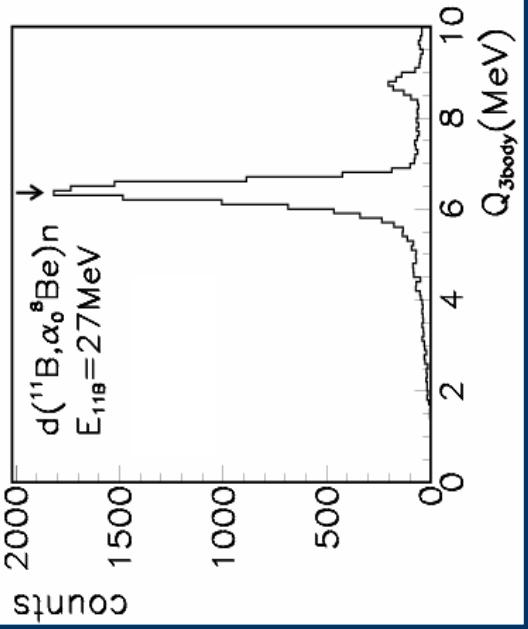
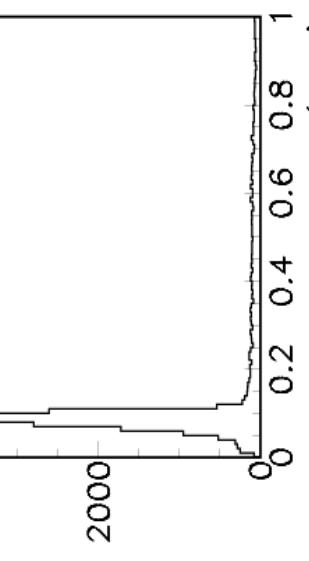
- The experiment was performed at LNS (Tandem & Camera2000) in two different runs on December 2002 and April 2007 ;
- $E_{\text{beam}}(^{11}\text{B})=27 \text{ MeV}$ & $I_{\text{beam}}(^{11}\text{B})=2-5 \text{ nA}$;
- Target thickness $CD_2 \sim 190 \mu\text{g}/\text{cm}^2$;
- Displacement of the detectors around the whole *QF-angular* range.

Determination of ${}^8\text{Be}$ and selection of the three-body channel ${}^2\text{H}({}^{11}\text{B}, \alpha_0 {}^8\text{Be})\text{n}$: 2007 Exp

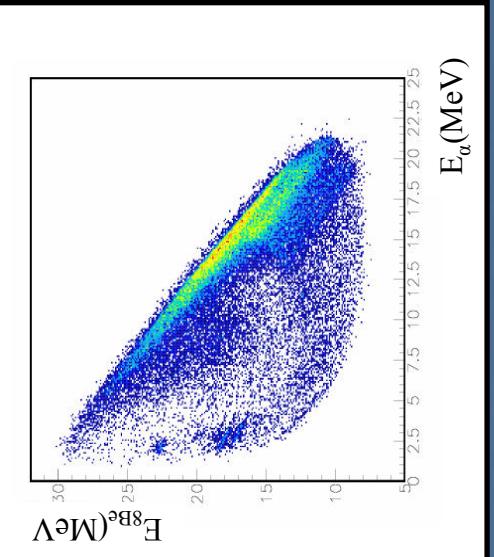
Alpha-decay of ${}^8\text{Be}$ (g.s.) in 2 alphas ($Q_{\text{dec}} \sim 90$ keV) detected in coincidence on DPSSD (Dual Position Sensitive Detector)



Off-line reconstruction
of relative energy for
alfa-decay of ${}^8\text{Be}(\text{g.s.})$



3-body kinematical locus: selection of channel



Q-value peak for the
3-body reaction
($Q_{\text{th}} = 6.36$ MeV)

Selection of the QF-mechanism:

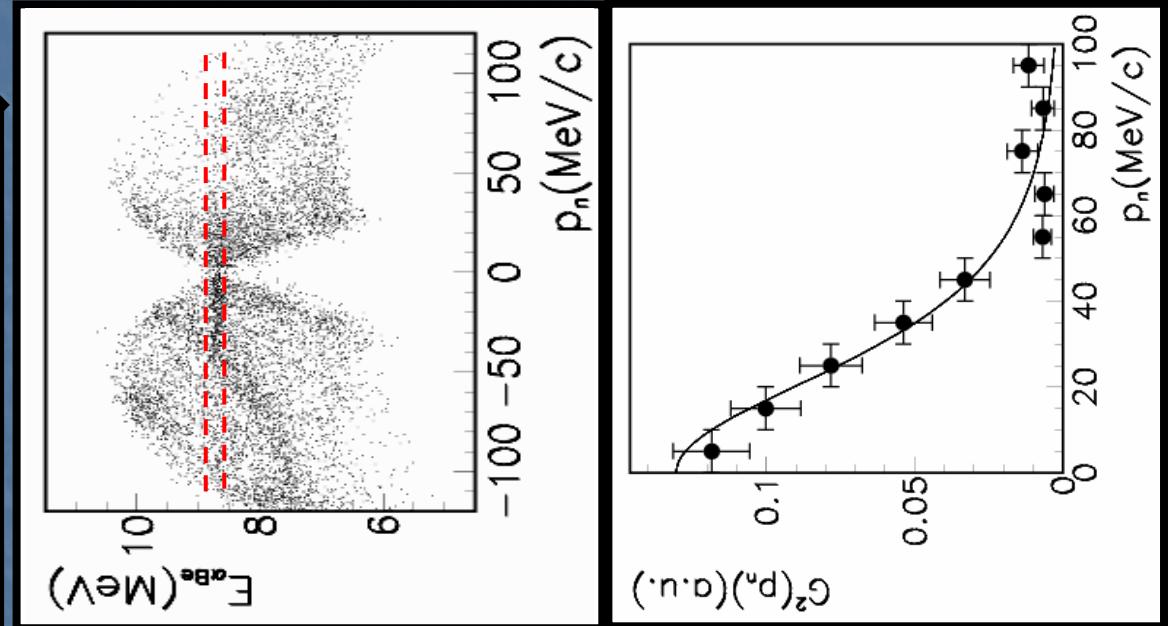
Experimental Momentum Distribution

✓ A further and necessary experimental check is to study the experimental neutron momentum distribution.

✓ Comparison between the experimental data (black points) and the theoretical Hulthén wave function ($a=0.2317 \text{ fm}^{-1}$, $b=1.202 \text{ fm}^{-1}$):

$$|\Phi(\vec{p}_n)|^2 = \frac{1}{2\pi} \sqrt{\frac{ab(a+b)}{(a-b)^2}} \left[\frac{1}{a^2 + p_n^2} - \frac{1}{b^2 + p_n^2} \right]$$

Necessary condition for the application of the Trojan Horse!!



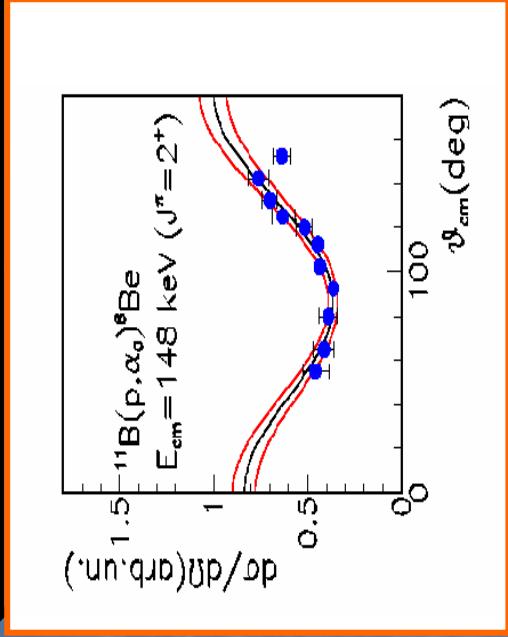
✓ QF-hypothesis satisfied



✓ THM application

$^{11}\text{B}(\text{p}, \alpha_0)^8\text{Be}$ angular distributions & S(E)-factor.

Comparison between Exp 2007 – Exp 2002



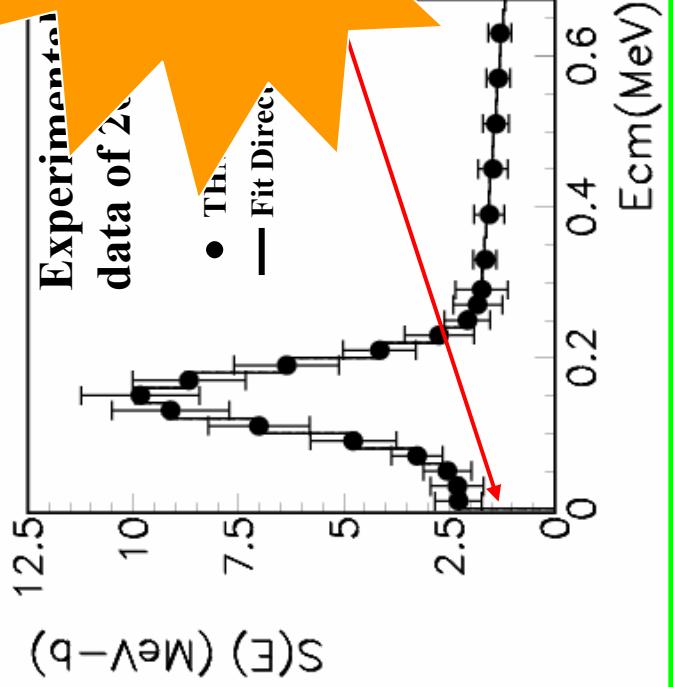
S(E)-extraction

$$S(E) = \frac{E * P_i}{\sigma^N(E)} * \exp(2\pi n)$$

APPROACH:

- ^{9}Be (SD) subtraction;
- Normalization between 400-600 keV's;
- Separation between resonant $l=1$ (16.106 MeV) and non-resonant $l=0$ contribution

Measurement
of
Gamow peak in
N-Gamma

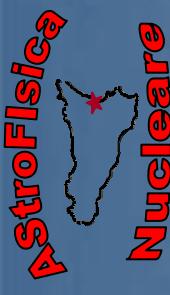


Experimental data of 2002

- THM Data
- Fit Direct



Final results, discussion & perspectives



■ S(E)-extraction

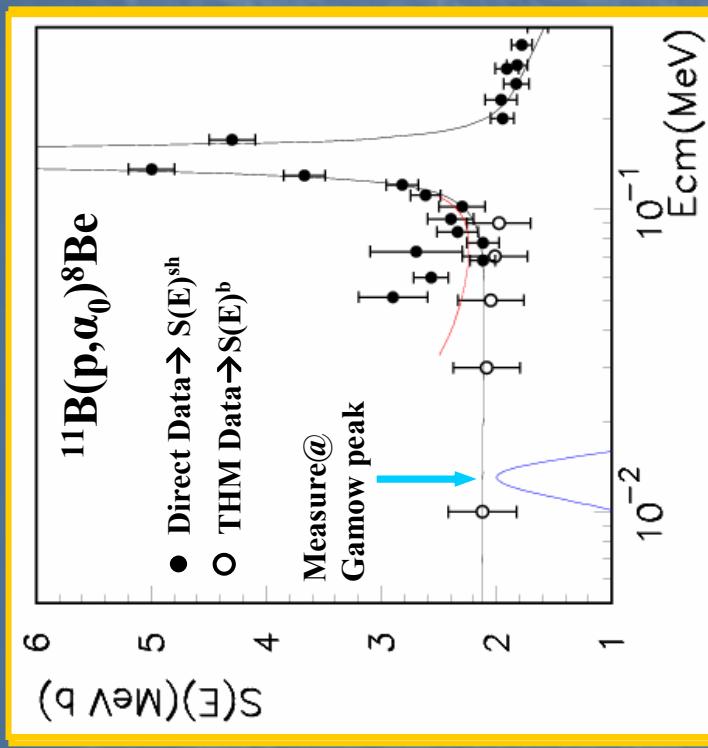
$$S(E) = E * P_l \sigma^N(E) * \exp(2\pi n)$$

The experimental Fit on THM data leads to

$$S(E)_{\text{THM NR}} = 2.14 - 1.85 * E + 0.89 * E^2 \text{ (MeV b)}$$



DATA	$S(E)$ VALUES	REFERENCES
THM 2007-2002	2.15 ± 0.28 MeV b	Present Work.
THM 2002 (new approach)	2.22 ± 0.3 MeV b	Lamia et al., N. Cimento, 2009
Extrapolations	2.10 ± 0.13 MeV b	Becker et al., Z.Ph.327, 341, 1987
PRC, 2004	0.41 ± 0.09 MeV b	Spitaleri et al., Phy.Rew.C, 2004



Final results, discussion & perspectives

► Non resonant reaction rate (“Nuclear Physics of Stars”, Iliadis C., 2007)

$$N_A \langle \sigma v \rangle = \frac{C_1}{T_g^{2/3}} e^{-C_2/T_g^{1/3}} \left(1 + C_3 T_g^{1/3} + C_4 T_g^{2/3} + C_5 T_g + C_6 T_g^{4/3} + C_7 T_g^{5/3} \right) \quad (\text{cm}^3 \text{mol}^{-1} \text{s}^{-1})$$

$$\begin{aligned} C_1 &= 7.8324 \times 10^9 \left(Z_0^2 Z_1^2 \frac{M_0 M_1}{M_0 + M_1} \right)^{1/6} \frac{\sqrt{M_0 + M_1}}{S(0)} \\ C_2 &= 4.2475 \left(Z_0^2 Z_1^2 \frac{M_0 M_1}{M_0 + M_1} \right)^{1/3} \\ C_3 &= 9.810 \times 10^{-2} \left(Z_0^2 Z_1^2 \frac{M_0 M_1}{M_0 + M_1} \right)^{-1/3} \\ C_4 &= 0.1220 \frac{S'(0)}{S(0)} \left(Z_0^2 Z_1^2 \frac{M_0 M_1}{M_0 + M_1} \right)^{1/3} \\ C_5 &= 8.377 \times 10^{-2} \frac{S'(0)}{S(0)} \\ C_6 &= 7.442 \times 10^{-3} \frac{S''(0)}{S(0)} \left(Z_0^2 Z_1^2 \frac{M_0 M_1}{M_0 + M_1} \right)^{2/3} \\ C_7 &= 1.299 \times 10^{-2} \frac{S''(0)}{S(0)} \left(Z_0^2 Z_1^2 \frac{M_0 M_1}{M_0 + M_1} \right)^{1/3} \end{aligned}$$

